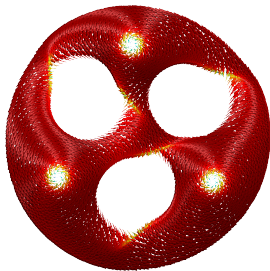




26th International Meshing Roundtable

Computing cross fields A PDE approach based on Ginzburg-Landau theory

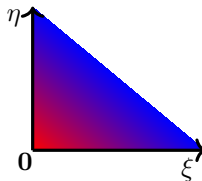
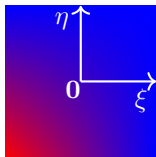
Pierre-Alexandre Beaufort, Jonathan Lambrechts, François Henrotte,
Christophe Geuzaine, Jean-François Remacle



Motivation

Meshing quadrangles

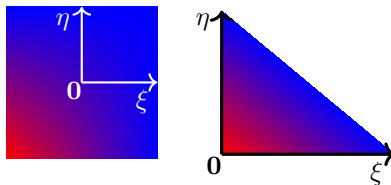
Meshing quadrangles for finite elements methods



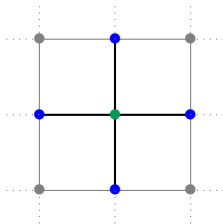
Motivation

Meshing quadrangles

Meshing quadrangles for finite elements methods



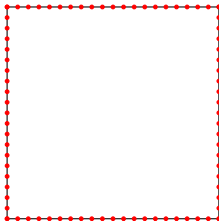
Quadrangle quality strongly depends on point locations



One way to spawn points

Frontal approach

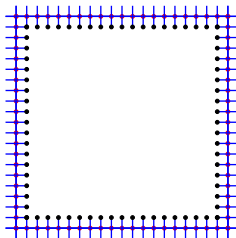
From boundaries,...



One way to spawn points

Frontal approach

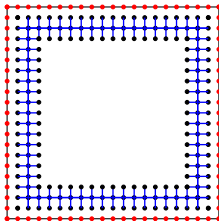
From boundaries,...



One way to spawn points

Frontal approach

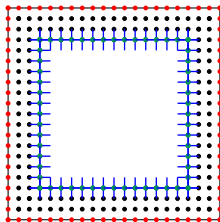
From boundaries,...



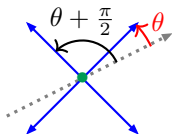
One way to spawn points

Frontal approach

From boundaries,...



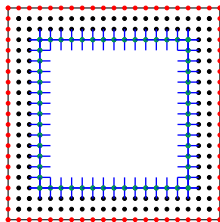
For all $\mathbf{x} \in M$, 4 preferred **orthonormal directions** are given



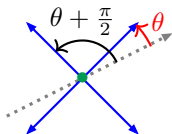
One way to spawn points

Frontal approach

From boundaries,...



For all $\mathbf{x} \in M$, 4 preferred **orthonormal directions** are given



\Rightarrow it defines a **cross field**

$$\theta \stackrel{?}{=} \theta + \frac{\pi}{2}k, \forall k \in \mathbb{Z}$$

Modeling

Cross fields



A cross field which relative angle is θ may be defined by

$$\langle \cos(4\theta); \sin(4\theta) \rangle$$

Modeling

Cross fields



A cross field which relative angle is θ may be defined by

$$\langle \cos(4\theta); \sin(4\theta) \rangle$$

which is suitable:

► **Uniqueness**

$$\cos\left(4\left[\theta + k\frac{\pi}{2}\right]\right) = \cos(4\theta), \quad \forall k \in \mathbb{Z}$$

► **Distance**

$$\begin{aligned} & \int_0^{2\pi} |\cos(4[\theta_i + \alpha]) - \cos(4[\theta_j + \alpha])|^2 d\alpha \\ &= \pi ((\cos(4\theta_i) - \cos(4\theta_j))^2 + (\sin(4\theta_i) - \sin(4\theta_j))^2) \end{aligned}$$

Complex analogy

Vector fields

Actually, a cross field consists of vector fields:

$$\underbrace{< \cos(4\theta); \sin(4\theta) >}_{u} \underbrace{\equiv}_{v} \underbrace{\exp(i \ 4\theta)}_{\exp(i \ \theta)^4} = u + i \ v$$

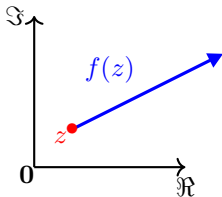
Complex analogy

Vector fields

Actually, a cross field consists of vector fields:

$$\langle \underbrace{\cos(4\theta)}_u; \underbrace{\sin(4\theta)}_v \rangle \equiv \underbrace{\exp(i \ 4\theta)}_{\exp(i \ \theta)^4} = u + i \ v$$

Two dimensional vector fields correspond to values of complex functions



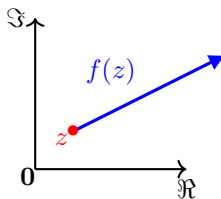
Complex analogy

Vector fields

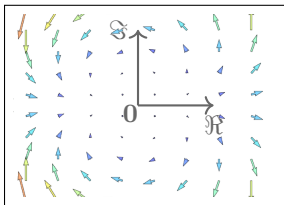
Actually, a cross field consists of vector fields:

$$\langle \underbrace{\cos(4\theta)}_u; \underbrace{\sin(4\theta)}_v \rangle \equiv \underbrace{\exp(i \ 4\theta)}_{\exp(i \ \theta)^4} = u + i \ v$$

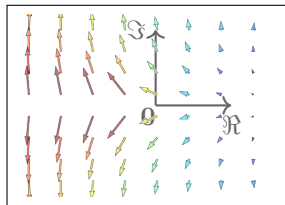
Two dimensional vector fields correspond to values of complex functions



$$f(z) = z^2$$



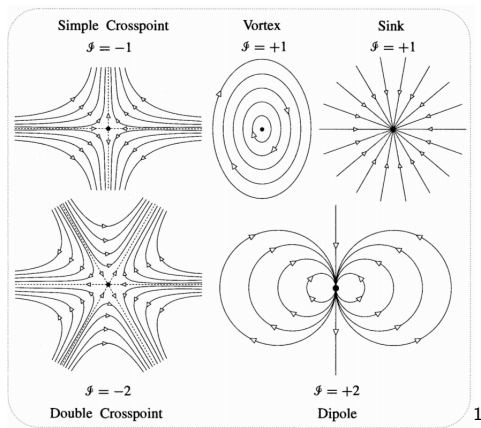
$$f(z) = \log(z)$$



Critical points

Indices

- ▶ Vector fields may have **critical point(s)** z^c , $f'(z^c) = 0$
- ▶ z^c has an **index** \mathcal{J} defined as winding number of points around



¹Figure from Tristan Needham's book, "Visual complex analysis"

Poincaré-Hopf theorem

Vector fields on closed surfaces

If a vector field on a smooth closed surface of genus g has only a finite number n of singular points s_j , then

$$\sum_{j=1}^n \mathfrak{I}(s_j) = 2(1 - g)$$

where $2(1 - g)$ equals the Euler characteristic of a closed surface

Computing cross fields

Criteria

How is a cross field built over a surface?

- ▶ Smooth cross fields for smooth directions
- ▶ Average boundary orientations within surface
- ▶ *Cross field should have unit norm almost everywhere**

Engineering approach

Energy formulation

Smooth out and average data from boundary conditions with Laplace

$$E(u; v) = \min_{u, v} \int_M |\nabla u|^2 + |\nabla v|^2 d\mathbf{x}$$

such that over ∂M : $u \equiv 1$ and $v \equiv 0$

Engineering approach

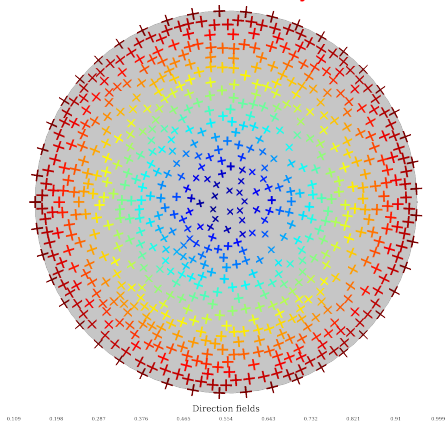
Energy formulation

Smooth out and average data from boundary conditions with Laplace

$$E(u; v) = \min_{u, v} \int_M |\nabla u|^2 + |\nabla v|^2 d\mathbf{x}$$

such that over ∂M : $u \equiv 1$ and $v \equiv 0$

But data vanishes far away boundaries



Engineering approach

Energy formulation

A penalty term is then added to foster unit norm cross fields

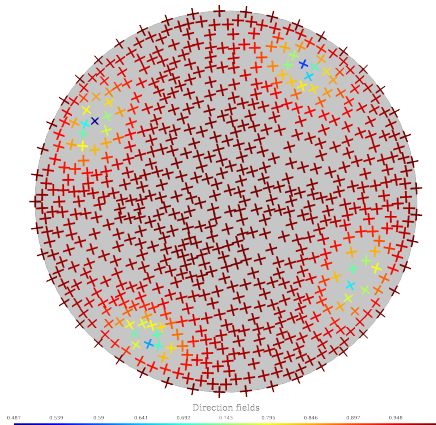
$$E(u; v) = \min_{u, v} \underbrace{\int_M |\nabla u|^2 + |\nabla v|^2 d\mathbf{x}}_{\text{Smoother term}} + \beta \underbrace{\int_M (u^2 + v^2 - 1)^2 d\mathbf{x}}_{\text{Penalty term}}$$

Engineering approach

Energy formulation

A penalty term is then added to foster unit norm cross fields

$$E(u; v) = \min_{u, v} \underbrace{\int_M |\nabla u|^2 + |\nabla v|^2 d\mathbf{x}}_{\text{Smoother term}} + \underbrace{\beta \int_M (u^2 + v^2 - 1)^2 d\mathbf{x}}_{\text{Penalty term}}$$



Ginzburg-Landau functional

Preliminaries

Ginzburg-Landau functional is

$$E_{\epsilon}(f) = \frac{1}{2} \int_M |\nabla f|^2 d\mathbf{x} + \frac{1}{4\epsilon^2} \int_M (|f|^2 - 1)^2 d\mathbf{x}$$

which is defined for maps $f \in H^1(M, \mathbb{C})$, i.e. $f \in H^1(M) : M \mapsto \mathbb{C}$

Ginzburg-Landau functional

Preliminaries

Ginzburg-Landau functional is

$$E_\epsilon(f) = \frac{1}{2} \int_M |\nabla f|^2 d\mathbf{x} + \frac{1}{4\epsilon^2} \int_M (|f|^2 - 1)^2 d\mathbf{x}$$

which is defined for maps $f \in H^1(M, \mathbb{C})$, i.e. $f \in H^1(M) : M \mapsto \mathbb{C}$

- ▶ ϵ is a characteristic length of M , called **coherence length**
- ▶ Let $H_g^1(M, \mathbb{C}) = \{f \in H^1(M, \mathbb{C}) : f = g \text{ on } \partial M\}$

Ginzburg-Landau functional

Preliminaries

Ginzburg-Landau functional is

$$E_\epsilon(f) = \frac{1}{2} \int_M |\nabla f|^2 d\mathbf{x} + \frac{1}{4\epsilon^2} \int_M (|f|^2 - 1)^2 d\mathbf{x}$$

which is defined for maps $f \in H^1(M, \mathbb{C})$, i.e. $f \in H^1(M) : M \mapsto \mathbb{C}$

- ▶ ϵ is a characteristic length of M , called **coherence length**
- ▶ Let $H_g^1(M, \mathbb{C}) = \{f \in H^1(M, \mathbb{C}) : f = g \text{ on } \partial M\}$

$$\min_{f \in H_g^1(M, \mathbb{C})} E_\epsilon(f)$$

Initial mapping

Topological requirements

Let $\mathcal{S}^1 = \{z \in \mathbb{C} : |z| = 1\}$.

$$\min_{f \in H_g^1(M, \mathcal{S}^1)} \int_M |\nabla f|^2 d\mathbf{x}$$

Solution corresponds to a smooth mapping between M and unit circle \mathcal{S}^1

Initial mapping

Topological requirements

Let $\mathcal{S}^1 = \{z \in \mathbb{C} : |z| = 1\}$.

$$\min_{f \in H_g^1(M, \mathcal{S}^1)} \int_M |\nabla f|^2 d\mathbf{x}$$

Solution corresponds to a smooth mapping between M and unit circle \mathcal{S}^1

Solution f_0 is unique and smooth
if and only if
 $\Im(g_{\partial M}) = 0$

Initial mapping

Topological requirements

Let $\mathcal{S}^1 = \{z \in \mathbb{C} : |z| = 1\}$.

$$\min_{f \in H_g^1(M, \mathcal{S}^1)} \int_M |\nabla f|^2 d\mathbf{x}$$

Solution corresponds to a smooth mapping between M and unit circle \mathcal{S}^1

Solution f_0 is unique and smooth
if and only if
 $\Im(g_{\partial M}) = 0$

Otherwise, $|\nabla f|^2$ is not bounded in some $\mathbf{x}^c \in M$

Relaxed mapping

Penalty term

If $\mathfrak{I}(g_{\partial M}) \neq 0 \implies H_g^1(M, \mathcal{S}^1) = \emptyset$

It means there does not exist admissible solution to minimization problem

Relaxed mapping

Penalty term

If $\mathfrak{I}(g_{\partial M}) \neq 0 \implies H_g^1(M, \mathcal{S}^1) = \emptyset$

It means there does not exist admissible solution to minimization problem

Hence, constraint is relaxed: it is implicitly enforced within formulation

$$\min_{f \in H_g^1(M, \mathbb{C})} \frac{1}{2} \int_M |\nabla f|^2 d\mathbf{x} + \frac{1}{4\epsilon^2} \int_M (|f|^2 - 1)^2 d\mathbf{x}$$

Relaxed mapping

Penalty term

If $\mathfrak{I}(g_{\partial M}) \neq 0 \implies H_g^1(M, \mathcal{S}^1) = \emptyset$

It means there does not exist admissible solution to minimization problem

Hence, constraint is relaxed: it is implicitly enforced within formulation

$$\min_{f \in H_g^1(M, \mathbb{C})} \frac{1}{2} \int_M |\nabla f|^2 d\mathbf{x} + \frac{1}{4\epsilon^2} \int_M (|f|^2 - 1)^2 d\mathbf{x}$$

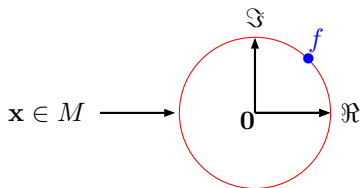
Asymptotic behavior of solution f_ϵ

$$\int_M |\nabla f_\epsilon|^2 d\mathbf{x} \xrightarrow{\epsilon \rightarrow 0} \infty$$

Link with directional fields

Ginzburg-Landau equation

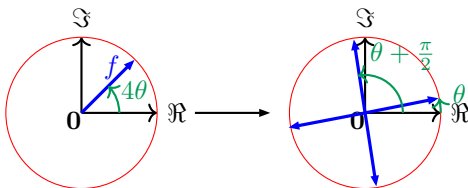
- Ginzburg-Landau functional maps $\mathbf{x} \in M$ over unit circle \mathcal{S}^1



Link with directional fields

Ginzburg-Landau equation

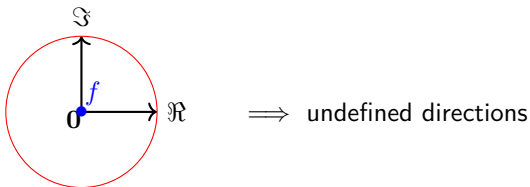
- Ginzburg-Landau functional maps $\mathbf{x} \in M$ over unit circle \mathcal{S}^1
- Mapping f may describe a directional field: $f = \exp(i 4\theta)$



Link with directional fields

Ginzburg-Landau equation

- ▶ Ginzburg-Landau functional maps $\mathbf{x} \in M$ over unit circle \mathcal{S}^1
- ▶ Mapping f may describe a directional field: $f = \exp(i 4\theta)$
- ▶ Asymptotic behavior of Ginzburg-Landau equation yields vector fields critical points



Critical points

Ginzburg-Landau functional

A critical point z^c has following contribution

$$\pi (\Im(z^c))^2 |\log(\epsilon)|$$

as ϵ tends to zero within Ginzburg-Landau functional E_ϵ

Critical points

Ginzburg-Landau functional

A critical point z^c has following contribution

$$\pi (\Im(z^c))^2 |\log(\epsilon)|$$

as ϵ tends to zero within Ginzburg-Landau functional E_ϵ

E_ϵ is minimum by minimizing absolute value of index of critical points

Critical points

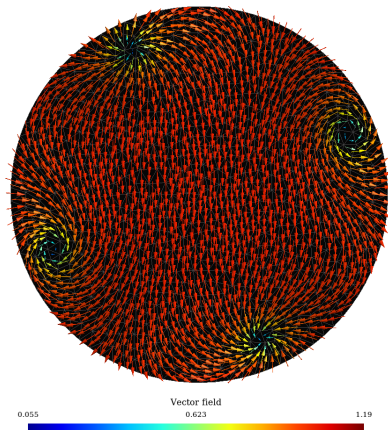
Ginzburg-Landau functional

A critical point z^c has following contribution

$$\pi (\Im(z^c))^2 |\log(\epsilon)|$$

as ϵ tends to zero within Ginzburg-Landau functional E_ϵ

E_ϵ is minimum by minimizing absolute value of index of critical points



Non unit norm cross field

Weak constraint

Since $|f|$ may be different of unity, it has to be redefined

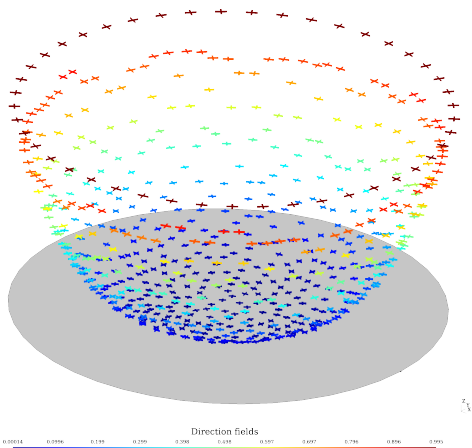
$$f := r \exp(i 4\theta)$$

Non unit norm cross field

Weak constraint

Since $|f|$ may be different of unity, it has to be redefined

$$f := r \exp(i 4\theta)$$



Non unit norm cross field

Weak constraint

Since $|f|$ may be different of unity, it has to be redefined

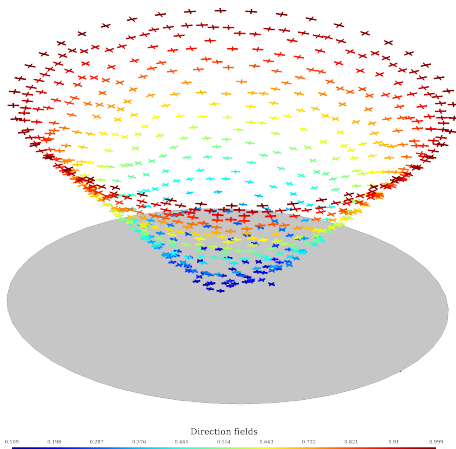
$$f := r^4 \exp(i \, 4\theta)$$

Non unit norm cross field

Weak constraint

Since $|f|$ may be different of unity, it has to be redefined

$$f := r^4 \exp(i 4\theta)$$



Directional fields

Interpretation

Directions of cross fields correspond to 4-th roots of vector fields expression

$$f(z) = z^4 = r^4 \exp(i 4\theta)$$

Directional fields

Interpretation

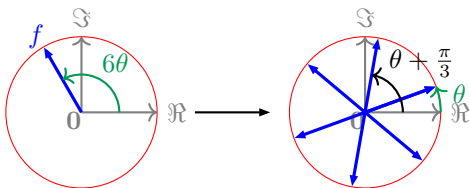
Directions of cross fields correspond to 4-th roots of vector fields expression

$$f(z) = z^4 = r^4 \exp(i 4\theta)$$

Directions of n-fields correspond to n-th roots of vector fields

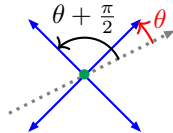
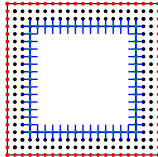
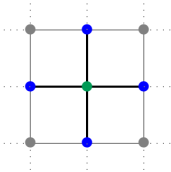
$$f(z) = z^n = r^n \exp(i n\theta)$$

Directional fields with 6 symmetries spawn vertices of equilateral triangles



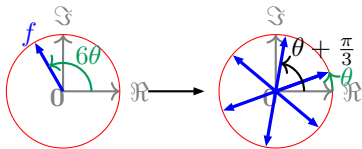
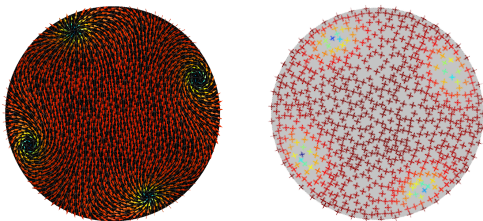
Conclusion

- Quadrangle quality may be ensured by a cross field



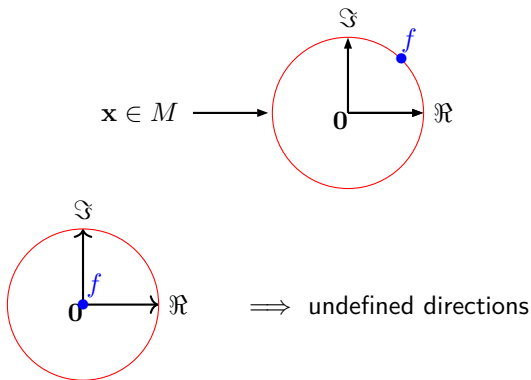
Conclusion

- ▶ Quadrangle quality may be ensured by a cross field
- ▶ A n -directional field corresponds to n -th root of a vector field



Conclusion

- ▶ Quadrangle quality may be ensured by a cross field
- ▶ A n -directional field corresponds to n -th root of a vector field
- ▶ $H_g^1(M, S^1)$ is too restrictive and unusable within FEM



Conclusion

- ▶ Quadrangle quality may be ensured by a cross field
- ▶ A n-directional field corresponds to n-th root of a vector field
- ▶ $H_g^1(M, S^1)$ is too restrictive and unusable within FEM
- ▶ Ginzburg-Landau functional is meaningful for flat or closed surfaces, and it is consistent for two-manifolds

$$\min_{f \in H_g^1(M, \mathbb{C})} \int_M |\nabla f|^2 d\mathbf{x} + \frac{1}{2\epsilon^2} \int_M (|f|^2 - 1)^2 d\mathbf{x}$$

$$E_\epsilon(\mathbf{x}^c) \approx \pi (\mathfrak{I}(\mathbf{x}^c))^2 |\log(\epsilon)|$$

For meshing examples, see Georgiadis 2017

For further details, see talk 6B.2 Jezdimirovic

Thank you for your attention!

Any questions?

This work is funded by **ARC WAVES** 15/19-03

[*] "Ginzburg-Landau vortices", F. Bethuel et al.

pierre-alexandre.beaufort@uclouvain.be